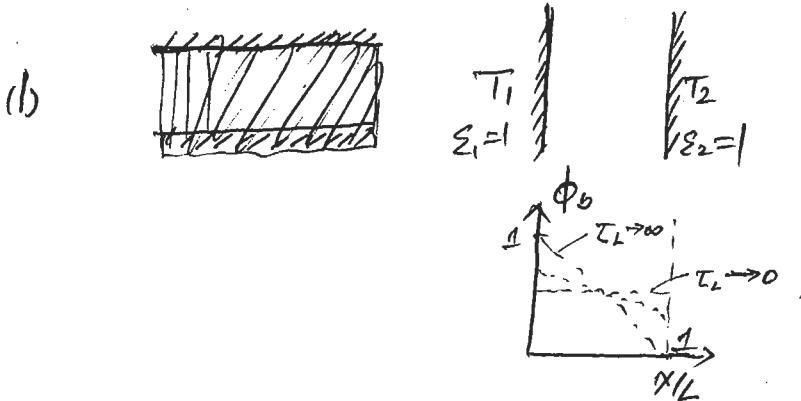


Review of last lecture



(2) approximate solutions

 $\tau_L \ll \phi$ optically thin $\tau_L \gg 1$ optically thick

$$\mu \frac{dI_\eta}{d\xi} = -I_\eta + I_{b\eta}$$

linear expansi.

but simply replace derivative term by $\frac{dI_\eta}{d\xi} \approx \frac{dI_{b\eta}}{d\xi}$

$$I_\eta = I_{b\eta} - \mu \frac{dI_{b\eta}}{d\xi}$$

$$\begin{aligned} q &= \int_0^\infty d\eta \int_{4\pi} I_\eta \cos\theta d\Omega \\ &= \int_0^\infty d\eta \int_0^{2\pi} d\varphi \int_0^\pi \left[I_{b\eta} - \mu \cos\theta \frac{dI_{b\eta}}{d\xi} \right] \sin\theta d\theta \\ &= -\frac{4\pi}{3} \int_0^\infty \frac{dI_{b\eta}}{d\xi} d\xi = -\frac{4\pi}{3} \frac{de_\eta}{d\xi} = -\frac{4}{3k_e} \frac{de_0}{dx} \\ &= -\frac{c}{3k_e} \frac{dU_b}{dx} = -\frac{c}{3k_e} \left(\frac{dU_b}{dT} \right) \frac{dT}{dx} = -\frac{1}{3} c v \frac{1}{k_e} \frac{dT}{dx} \end{aligned}$$

Approximate Solutions:

$\tau_L \ll 1$ optically thin

~~$\frac{dI_\eta}{dz} = -I_\eta + I_{b\eta}$~~ $q = (J_1 - J_2)(1 - \tau_L)$
 \uparrow
 linearly attenuated

$\tau_L \gg 1$ optically thick:

$$\mu \frac{dI_\eta}{dz} = -I_\eta + I_{b\eta}$$

linear expansion at $I_{b\eta}$ $I_\eta = I_{b\eta} + \beta \frac{dI_{b\eta}}{dz} + \dots$

1st order solution: $\beta = \bar{\mu}$

$$I_\eta = I_{b\eta} + \bar{\mu} \frac{dI_{b\eta}}{dz}$$

$$q = \int_0^\infty \int_0^{2\pi} d\eta \int_0^\pi \sin\theta d\theta \left[\int_0^\infty I_\eta \mu d\mu - \int_0^\infty I_{b\eta} \mu d\mu \right]$$

$$= \int_0^\infty d\eta \int_{4\pi} I_\eta \mu d\Omega$$

$$= \int_0^\infty d\eta \int_0^{2\pi} d\phi \int_0^\pi \sin\theta d\theta \left[I_{b\eta} + \bar{\mu} \frac{dI_{b\eta}}{dz} \right] \mu d(-\mu)$$

$$= +2\pi \int_0^\pi \mu^2 \left(\int_0^\infty \frac{dI_{b\eta}}{dz} d\eta \right) d\mu$$

energy density \uparrow

$$= -\frac{4\pi}{3} \frac{dI_b}{dz} = -\frac{16\pi}{3k_e} \frac{d(4\pi I_b)}{dz} = -\frac{c k_e}{3k_e} \frac{d(u)}{dz} = -\frac{c}{3k_e} \frac{d(u)}{dz}$$

$$= -\frac{4\pi}{3k_e} \frac{d(\sigma T^4)}{dz} = -\frac{c}{3k_e} \frac{dT}{dz}$$

General

$$q = - \frac{4\pi}{3k_e} \nabla I_b \quad \text{--- Rosseland diffusion approximation}$$

Boundary Conditions:

Diffusion approximation right to the wall



$$\text{Int. flux leaving wall: } I^+ = \epsilon I_{bw} + I^- (1-\epsilon)$$

$$q = \int_0^{2\pi} d\varphi \left[\int_0^1 I^+ \mu d\mu - \int_0^1 I^- (-\mu) d\mu \right]$$

$$= 2\pi \left[\epsilon \cdot \frac{1}{2} I_{bw} + \epsilon \int_0^1 I^- (-\mu) d\mu \right]$$

$$I^- (\mu) = I_b^{(0)} - \mu \frac{dI_b}{dz}$$

$$q = 2\pi \epsilon \left[\frac{1}{2} I_{bw} - \frac{1}{3} \frac{dI_b}{dz} \right]$$

$$q = \epsilon \left[E_{bw} - E_b(0) \right] + \frac{2q}{3}$$

$$q = \frac{\epsilon [E_{bw} - E_b(0)]}{\frac{1}{2} - \frac{1}{3}}$$

$$q(L) = \frac{E_b(L) - E_{bw2}}{\frac{1}{2} - \frac{1}{3}}$$

If radiative equilibrium $q = \text{constant}$

$$E_b(0) - E_b(L) = \frac{3k_e L}{4} q$$

$$\Rightarrow q = \frac{E_{bw1} - E_{bw2}}{\frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} + \frac{3}{4} \tau L}$$

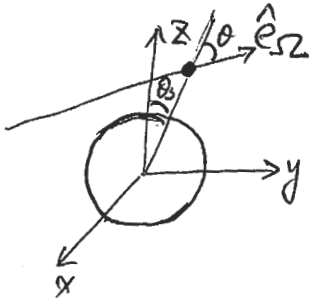
Comments
 (1) higher order expansion
 (2)

Spherical & Cylindrical Coordinates:

2014

$$\frac{dI_\eta}{ds} = \hat{e}_\Omega \cdot \nabla_r I_\eta$$

$$\frac{1}{ke} \hat{e}_\Omega \cdot \nabla_r I_\eta = -I_\eta + \underbrace{(1-\cos\theta) I_{b\eta}}_S + \frac{10\eta}{4\pi} \int \Phi(\Omega' \rightarrow \Omega) I(\Omega') d\Omega'$$

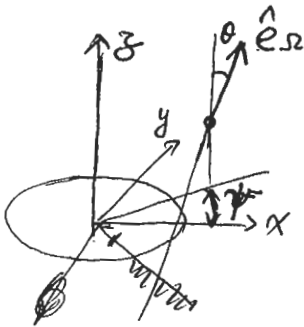


$$\mu = \cos\theta$$

$$\mu \frac{dI_\eta}{dz} + \frac{1-\mu^2}{z} \frac{dI_\eta}{d\mu} = S_\eta - I_\eta$$

↑
local

Cylindrical Coordinate:



$$\sin\theta \left[\cos\psi \frac{\partial I_\eta}{\partial z} - \frac{\sin\psi}{z} \frac{\partial I_\eta}{\partial \psi} \right] = S_\eta - I_\eta$$

Diffusion approximation: $\frac{\partial I_\eta}{\partial z}, \frac{\partial I_\eta}{\partial \psi} \rightarrow \frac{\partial I_{b\eta}}{\partial z}, \frac{\partial I_{b\eta}}{\partial \psi}$

z/2

Numerical Solution: Discrete Ordinate Method

$$z = \int_{4\pi} I \cos\theta d\Omega$$

$$2\pi \left[\int_0^1 I^+ \mu d\mu - \int_0^1 I^-(-\mu) \mu d\mu \right]$$

↑
axisymmetric