

2.003J/1.053J Dynamics and Control I, Spring 2007
 Professor Thomas Peacock
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Lecture 9

2D Motion of Rigid Bodies: Kinetics, Poolball Example

Kinetics of Rigid Bodies

Angular Momentum Principle for a Rigid Body

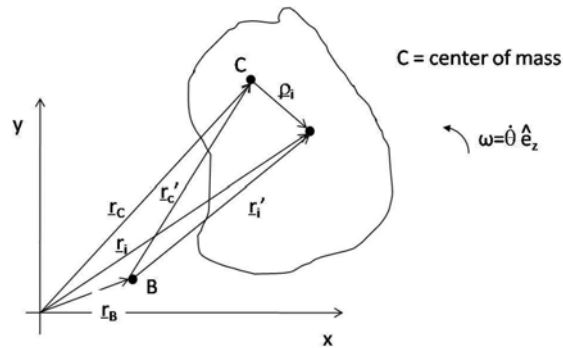


Figure 1: Rigid Body rotating with angular velocity ω . Figure by MIT OCW.

$$\underline{H}_B = \sum_i \underline{r}'_i \times m_i (\underline{v}_c + \underline{\omega} \times \underline{\rho}'_i)$$

After some steps (see Lecture 8):

$$\underline{H}_B = \underline{r}'_c \times \underline{P} + \sum_i m_i \underline{\rho}'_i \times \underline{\omega} \times \underline{\rho}'_i$$

We now use:

$$\underline{a} \times \underline{b} \times \underline{c} = (\underline{a} \cdot \underline{c})\underline{b} - (\underline{a} \cdot \underline{b})\underline{c}$$

$$\begin{aligned} \underline{\rho}'_i \times \underline{\omega} \times \underline{\rho}'_i &= \rho_i^2 \underline{\omega} - (\underline{\omega} \cdot \underline{\rho}'_i)\underline{\rho}'_i \\ &= \rho_i^2 \underline{\omega} \end{aligned}$$

For 2-D motion, $\underline{\omega} \cdot \underline{\rho}_i = 0$ because the vectors are \perp . For 3-D, this term does not have to be 0.

$$\begin{aligned}\underline{H}_B &= \underline{r}'_c \times \underline{P} + \sum_i m_i \rho_i^2 \underline{\omega} \\ &= \underline{r}'_c \times \underline{P} + I_c \underline{\omega}\end{aligned}$$

I_c : Moment of Inertia. $I_c = \sum_i m_i \rho_i^2$ (Intrinsic Property of Rigid Body)

Example:

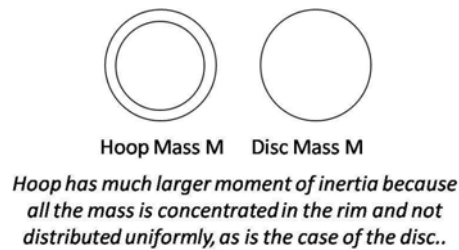


Figure 2: Hoop and Disc, both with mass M . Figure by MIT OCW.

$$\underline{H}_B = \underline{r}'_c \times \underline{P} + I_c \underline{\omega}$$

If one takes angular momentum about the center of mass:

$$\underline{H}_c = I_c \underline{\omega}$$

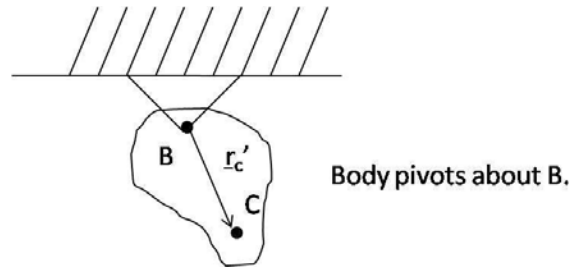
(Angular Momentum about B) = (Angular Momentum about C) + (Moment of Linear Momentum about B)

Therefore:

$$\underline{H}_B = \underline{H}_c + \underline{r}'_c \times \underline{P}$$

Special Case of Fixed Axis of Rotation about B

i.e. $\underline{v}_c = \underline{v}_B + \underline{\omega} \times \underline{r}'_c$

Figure 3: Rigid body which pivots about B . Figure by MIT OCW.

$$\begin{aligned}\underline{H}_B &= \underline{H}_C + \underline{r}'_c \times m(\underline{\omega} \times \underline{r}'_c) \\ &= \underline{H}_C + mr'_C{}^2 \underline{\omega} \\ &= (I_C + mr'_C{}^2) \underline{\omega} = I_B \underline{\omega}\end{aligned}$$

$$I_B = I_C + mr'_C{}^2 \quad \text{Parallel Axis Theorem}$$

$$\text{Only do this if the } \underline{v}_B = 0 \text{ and } \underline{v}_C = (\underline{\omega} \times \underline{r}'_C)$$

Finally:

$$\tau_B^{ext} = \frac{d}{dt} H_B + \underline{v}_B \times \underline{P} \quad (1)$$

$$\underline{H}_B = \underline{H}_C + \underline{r}'_C \times \underline{P} \quad (2)$$

$$\underline{H}_C = I_C \underline{\omega} \quad (3)$$

$$I_C = \sum m_i \rho_i^2 \quad (4)$$

Equations (1) to (4) are always true.

Special Cases

$$1. B = C \Rightarrow \underline{r}'_C = 0; \underline{v}_B \parallel \underline{P}$$

Start by thinking about motion around center of mass.

$$\tau_B^{ext} = \frac{d}{dt} \underline{H}_C \text{ and } \underline{H}_C = I_C \underline{\omega}$$

2. B is a stationary point and fixed in the body.

$$\tau_B^{ext} = \frac{d}{dt} \underline{H}_B \text{ and } \underline{H}_B = I_B \underline{\omega} \text{ where } I_B = I_C + mr_C'^2$$

What do we need to do still?

Calculating moments of inertia \Rightarrow Recitation 5

Work-Energy Principle

Cue hitting a pool ball

A pool ball of radius R and mass M is at rest on a horizontal table. It is set in motion by a sharp horizontal impulse \underline{J} provided by the cue. Determine the height above the ball's center that the cue should strike so that the subsequent motion is rolling without slipping.

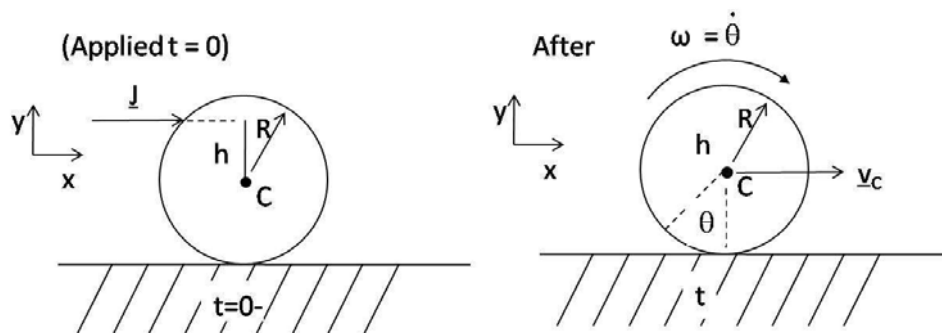


Figure 4: Cue ball diagram. Diagram shows cue ball when force is first applied and after being hit. Figure by MIT OCW.

Hit below h : Backspin

Hit above h : Top spin, carry on shot

Kinematics: Geometry with no forces

Horizontal Table: $y_C = \text{constant} = R$

Rolling without slipping: $v_C = \omega R$ (or $x_C = R\theta$)

1 Degree of Freedom. (Use x_C or θ).

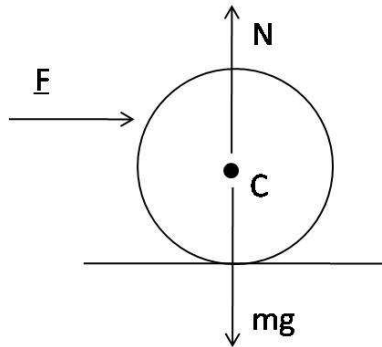
Kinetics: Free Body Diagrams

Figure 5: Free Body Diagram of Cue Ball. Figure by MIT OCW.

Impulse force that provides impulse \underline{J}

$$\underline{J} = \int_{0^-}^{0^+} \underline{F} dt = \int_{0^-}^{0^+} \underline{J}\delta(t) dt \text{ i.e. } \underline{F} = \underline{J}\delta(t)$$

(i) Linear Momentum Principle

$$\underline{F}^{ext} = \frac{d}{dt} \underline{P}$$

y-direction: C always at same height. $N = mg$ so no vertical motion of C .

x-direction: $F = J\delta(t) = \frac{d}{dt} Mv_C$.

Integrate both sides

$$\int_{0^-}^{0^+} F dt \int_{0^-}^{0^+} J\delta(t) dt = \int_{0^-}^{0^+} \frac{d}{dt} Mv_C dt$$

$$J = Mv_C(0^+) - Mv_C(0^-)$$

J : Momentum Imparted

$$\boxed{J = Mv_C(0^+)}$$

(5)

Angular Momentum Principle About C

Taking momentum about C simplifies equations

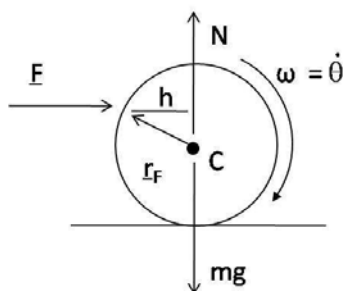


Figure 6: Angular Momentum Principle about C applied to Cue Ball. Figure by MIT OCW.

$$\begin{aligned}\underline{\tau}_C^{ext} &= \frac{d}{dt}\underline{H}_C \text{ and } \underline{H}_C = I_C\underline{\omega} \\ \underline{r}_F \times \underline{F} &= \frac{d}{dt}I_C\underline{\omega} \\ -Fh\hat{e}_z &= -I_C\frac{d\underline{\omega}}{dt}\hat{e}_z \\ \int_{0^-}^{0^+} \underline{F}h dt &= \int_{0^-}^{0^+} I_C\frac{d\underline{\omega}}{dt} dt \\ \int_{0^-}^{0^+} J\delta(t) dt &= I_C\underline{\omega}(0^+) - I_C\underline{\omega}(0^-)\end{aligned}$$

$I_C\underline{\omega}(0^-) = 0$ because $\omega(0^-) = 0$.

$$\boxed{Jh = I_C\underline{\omega}(0^+)}$$

Impulsive torque about center of mass = Change of angular momentum caused by the torque

Satisfying Constraints

If there is no slip, one needs $\omega(0^+)R = v_C(0^+)$

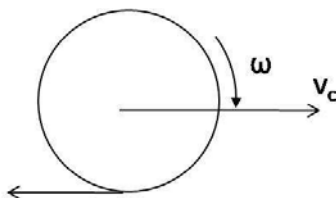


Figure 7: Diagram of Cue Ball moving. This diagram demonstrates how to satisfy geometric constraints of movement. Figure by MIT OCW.

$$J = \frac{I_C v_C(0^+)}{h R} \quad (6)$$

Can eliminate J from Equation 5 and Equation 6.

$$M v_C(0^+) = \frac{I_C v_C(0^+)}{h R} \Rightarrow \boxed{h = \frac{I_C}{mR}}$$

For a sphere:

$$I_C = \frac{2}{5} m R^2 \Rightarrow \boxed{h = \frac{2}{5} R}$$

h : Independent of mass of sphere. Independent of force applied.