

MIT OpenCourseWare
<http://ocw.mit.edu>

6.453 Quantum Optical Communication
Spring 2009

For information about citing these materials or our Terms of Use, visit: <http://ocw.mit.edu/terms>.



November 18, 2008

6.453 *Quantum Optical Communication* Lecture 18

Jeffrey H. Shapiro

Optical and Quantum Communications Group

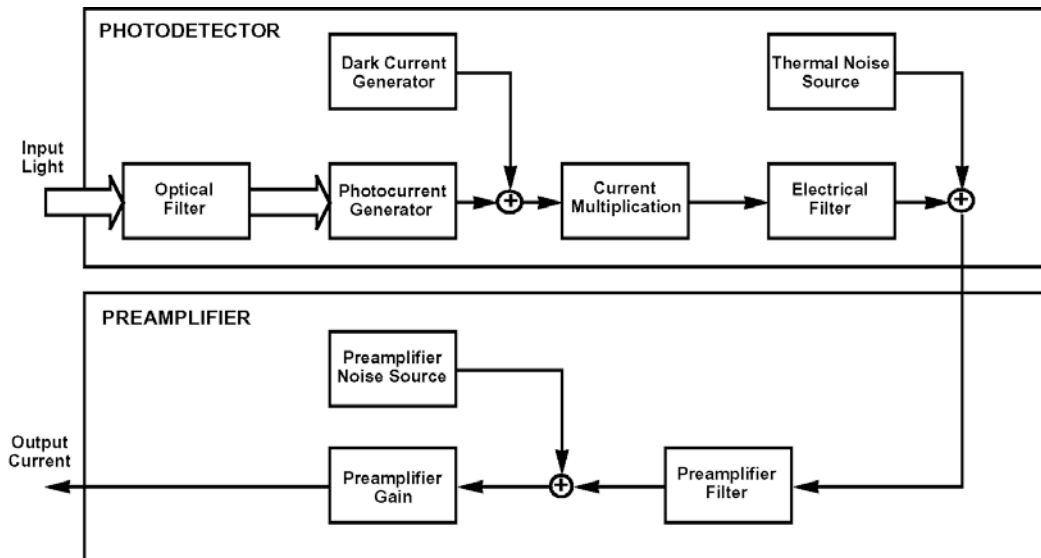
RESEARCH LABORATORY OF ELECTRONICS
Massachusetts Institute of Technology

www.rle.mit.edu/qoptics

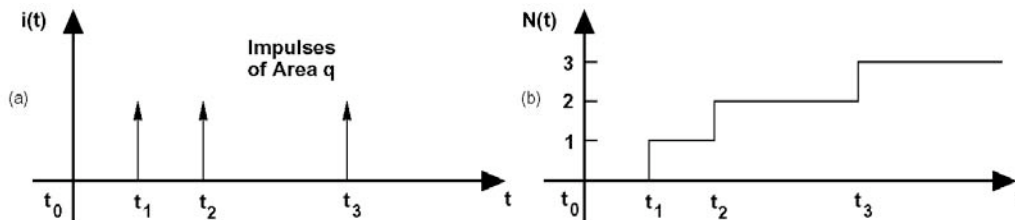
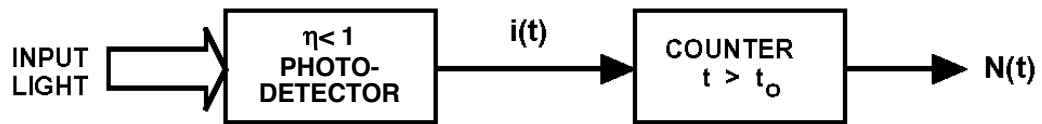
6.453 *Quantum Optical Communication* - Lecture 18

- Announcements
 - Pick up random processes notes, lecture notes, slides
- Continuous-Time Photodetection
 - Semiclassical theory — Poisson shot noise
 - Quantum theory — photon-flux operator measurement
 - Direct-detection signatures of non-classical light

Real Photodetection Systems



"Ideal" Photodetection System: Efficiency < 1



Classical Field versus Quantum Field Operator

- For Semiclassical Photodetection of Narrowband Light
 - Illumination is a classical photon-units positive-frequency field:

$$E(t)e^{-j\omega_o t}$$

- Short-Time Average Power on Detector:

$$P(t) = \hbar\omega_o |E(t)|^2$$

- For Quantum Photodetection of Narrowband Light
 - Illumination is a photon-units positive-frequency field operator:

$$\hat{E}(t)e^{-j\omega_o t}$$

- Only non-vacuum frequency components are within $\Delta\omega \ll \omega_o$

Semiclassical versus Quantum Photodetection

- Semiclassical Theory: Given $\{ P(t) : 0 \leq t \leq T \}$
 - $\{ N(t) : 0 \leq t \leq T \}$ is an inhomogeneous Poisson Counting Process
 - Rate function $\lambda(t) \equiv \eta P(t) / \hbar\omega_o$

- Quantum Theory:

$$\hat{E}'(t) \equiv \sqrt{\eta} \hat{E}(t) + \sqrt{1-\eta} \hat{E}_\eta(t)$$

$$i(t) \leftrightarrow \hat{i}(t) \equiv q \hat{E}'^\dagger(t) \hat{E}'(t)$$

$$N(t) \leftrightarrow \frac{1}{q} \int_0^t d\tau \hat{i}(\tau), \quad \text{for } 0 \leq t \leq T$$

Inhomogeneous Poisson Counting Process (IPCP)

- Definition: $\{N(t) : t \geq 0\}$ is an IPCP with rate $\lambda(t)$
 - Starts counting at zero: $N(0) = 0$
 - Has statistically independent increments
 - Increments are Poisson distributed

$$\Pr(N(t) - N(u) = n) = \frac{\left(\int_u^t d\tau \lambda(\tau)\right)^n \exp\left(-\int_u^t d\tau \lambda(\tau)\right)}{n!}$$

Mean and Covariance: Deterministic Illumination

- Semiclassical Photocount and Photocurrent Mean Functions:

$$\langle N(t) \rangle = \int_0^t d\tau \frac{\eta P(\tau)}{\hbar\omega_o} \quad \text{and} \quad \langle i(t) \rangle = q \frac{\eta P(t)}{\hbar\omega_o}$$

- Semiclassical Covariance Functions:

$$\begin{aligned} \langle \Delta N(t) \Delta N(u) \rangle &= \langle N(\min(t, u)) \rangle \\ \langle \Delta i(t) \Delta i(u) \rangle &= q \langle i(t) \rangle \delta(t - u) \end{aligned}$$

Mean and Covariance: Random Illumination

- Semiclassical Photocount and Photocurrent Mean Functions:

$$\langle N(t) \rangle = \int_0^t d\tau \frac{\eta \langle P(\tau) \rangle}{\hbar\omega_o}$$

$$\langle i(t) \rangle = q \frac{\eta \langle P(\tau) \rangle}{\hbar\omega_o}$$

- Semiclassical Covariance Functions:

$$\langle \Delta N(t) \Delta N(u) \rangle = \underbrace{\langle N(\min(t, u)) \rangle}_{\text{shot noise}} + \underbrace{\int_0^t d\tau \int_0^u d\tau' \frac{\eta^2 \langle \Delta P(\tau) \Delta P(\tau') \rangle}{(\hbar\omega_o)^2}}_{\text{excess noise}}$$

$$\langle \Delta i(t) \Delta i(u) \rangle = \underbrace{q \langle i(t) \rangle \delta(t - u)}_{\text{shot noise}} + \underbrace{q^2 \frac{\eta^2 \langle \Delta P(t) \Delta P(u) \rangle}{(\hbar\omega_o)^2}}_{\text{excess noise}}$$

Mean and Covariance Functions: Quantum Case

- Quantum Photocount and Photocurrent Mean Functions:

$$\langle N(t) \rangle = \int_0^t d\tau \eta \langle \hat{E}^\dagger(\tau) \hat{E}(\tau) \rangle \quad \text{and} \quad \langle i(t) \rangle = q\eta \langle \hat{E}^\dagger(t) \hat{E}(t) \rangle$$

- Quantum Covariance Functions:

$$\langle \Delta N(t) \Delta N(u) \rangle = \langle N(\min(t, u)) \rangle + \int_0^t d\tau \int_0^u d\tau' \eta^2 \left[\langle \hat{E}^\dagger(\tau) \hat{E}^\dagger(\tau') \hat{E}(\tau) \hat{E}(\tau') \rangle - \langle \hat{E}^\dagger(\tau) \hat{E}(\tau) \rangle \langle \hat{E}^\dagger(\tau') \hat{E}(\tau') \rangle \right]$$

$$\langle \Delta i(t) \Delta i(u) \rangle = q \langle i(t) \rangle \delta(t - u) + q^2 \eta^2 \left[\langle \hat{E}^\dagger(t) \hat{E}^\dagger(u) \hat{E}(t) \hat{E}(u) \rangle - \langle \hat{E}^\dagger(t) \hat{E}(t) \rangle \langle \hat{E}^\dagger(u) \hat{E}(u) \rangle \right]$$

Direct-Detection Signatures of Non-Classical Light

- Semiclassical Theory is Quantitatively Correct:
 - For coherent-state inputs $|E(t)\rangle$
 - For inputs that are classically-random mixtures of coherent states

- Sub-Poissonian Photon Counting:

- Semiclassical theory:

$$\langle \Delta N^2(t) \rangle \geq \langle N(t) \rangle$$

- Quantum theory:

$$\langle \Delta N^2(t) \rangle \geq 0$$

- Non-classical signature:

$$\langle \Delta N^2(t) \rangle < \langle N(t) \rangle$$

Direct-Detection Signatures of Non-Classical Light

- Photocurrent Noise Spectral Density for CW Sources:

$$\mathcal{S}_{ii}(\omega) \equiv \int_{-\infty}^{\infty} d\tau \langle \Delta i(t + \tau) \Delta i(t) \rangle e^{-j\omega\tau}$$

- Semiclassical Theory:

$$\mathcal{S}_{ii}(\omega) = q\langle i \rangle + q^2 \frac{\eta^2 \mathcal{S}_{PP}(\omega)}{(\hbar\omega_o)^2} \geq q\langle i \rangle$$

- Quantum Theory:

$$\mathcal{S}_{ii}(\omega) \geq 0$$

- Sub-Shot-Noise Non-Classical Signature:

$$\mathcal{S}_{ii}(\omega) < q\langle i \rangle$$

Coming Attractions: Lecture 19

- Lecture 19:
Continuous-Time Photodetection
 - Noise spectral densities in direct detection
 - Semiclassical theory of coherent detection
 - Quantum theory of coherent detection
 - Coherent-detection signatures of non-classical light